## Physics 735: Exam II (Take Home - Makeup)

March 24, 2009 Due March 25, 2009, 17:00 (5 pm)

## 1. New Diagonal Form of the Dirac Equation

We can introduce a Foldy-Wouthuysen transformation which will completely eliminate the lower component from the free positive energy solutions and the upper components from the free negative energy solutions.

The advantage of such a representation is that it allows us to regard the mixing of upper and lower components as a dynamical consequence of the interaction, whereas the free Dirac equation is fully diagonalized.

A unitary transformation which accomplishes this is

$$U = \sqrt{\frac{E+m}{2E}} \begin{pmatrix} 1 & \frac{\vec{\sigma} \cdot \vec{p}}{E+m} \\ \frac{-\vec{\sigma} \cdot \vec{p}}{E+m} & 1 \end{pmatrix}$$
(1)

- (a) [4 pts] Show by direct computation that U is unitary, i.e.  $U^{\dagger}U = \mathbf{1}$ .
- (b) 2 pts Show that

$$U u(\vec{p}, s) = \sqrt{2E} \begin{pmatrix} \chi^{(s)} \\ 0 \end{pmatrix}$$
$$U v(-\vec{p}, -s) = \sqrt{2E} \begin{pmatrix} 0 \\ -i\sigma_2 \chi^{(-s)} \end{pmatrix}.$$
 (2)

where

$$u(\vec{p},s) = \sqrt{E_p + m} \begin{pmatrix} 1\\ \frac{\vec{\sigma}\cdot\vec{p}}{E_p + m} \end{pmatrix} \chi^{(s)}$$
$$v(\vec{p},s) = \sqrt{E_p + m} \begin{pmatrix} \frac{\vec{\sigma}\cdot\vec{p}}{E_p + m}\\ 1 \end{pmatrix} \begin{bmatrix} -i\sigma_2\chi^{(s)} \end{bmatrix}.$$
(3)

Here  $\chi^{(s)}$  is the usual Pauli spinor.

(c) [5 pts] Show that the transformed free Hamiltonian,  $H_0 = \vec{\alpha} \cdot \vec{p} + \beta m$  is given by

$$UH_0 U^{\dagger} = \beta E = \begin{pmatrix} E & 0 & 0 & 0 \\ 0 & E & 0 & 0 \\ 0 & 0 & -E & 0 \\ 0 & 0 & 0 & -E \end{pmatrix}.$$
 (4)

In this form the free Dirac equation takes a form similar to the 2-component Klein-Gordon equation.

(d) [15 pts] Show that in this representation the Dirac equation with electromagnetic interactions (given usually with  $H = \vec{\alpha} \cdot (\vec{p} - \vec{V} + \beta m + V^0)$  can be written

$$i\frac{\partial}{\partial t}\psi = \beta E\psi + (H_1 + \beta H_2 + \vec{\sigma} \cdot \vec{H}_3 + \beta \vec{\sigma} \cdot \vec{H}_4 + \vec{\alpha} \cdot \vec{H}_5 + i\beta \vec{\alpha} \cdot \vec{H}_6 + \gamma^5 H_7)\psi$$
(5)

Proceed by introducing

$$\vec{D} = \frac{\vec{p}}{E+m} = \frac{-i\nabla}{E_{\nabla}+m}$$

$$\theta = \sqrt{\frac{E+m}{2E}}$$
(6)

and show that the functions  $H_i(i = 1, 7)$  are given by

$$H_{1} = \theta(V^{0} + D_{j}V^{0}D_{j})\theta$$

$$H_{2} = -\theta(\vec{V} \cdot \vec{D} + \vec{D} \cdot \vec{V})\theta$$

$$H_{3}^{k} = \theta(i\epsilon_{ijk}D_{i}V^{0}D_{j})\theta$$

$$\vec{H}_{4} = -\theta(i\vec{V} \times \vec{D}) + i(\vec{D} \times \vec{V}))\theta$$

$$\vec{H}_{5} = \theta(-\vec{V} + (\vec{D} \cdot \vec{V})\vec{D} + \vec{D}(\vec{V} \cdot \vec{D}) - \sum_{j}D_{j}\vec{V}D_{j})\theta$$

$$\vec{H}_{6} = \theta(iV^{0}\vec{D} - i\vec{D}V^{0})\theta$$

$$H_{7} = \theta(i\epsilon_{ijk}D_{i}V_{j}D_{k})\theta$$
(7)

## 2. Massless Dirac particle in a 1D square well

A massless spin-1/2 particle moves in a one-dimensional square well potential of depth  $V_0 \leq 0$ and a width *a*. Assume that the potential is a scalar potential (i.e. behaves as  $\beta V_0$ ). Divide the real axis into three domains:

$$I: \quad z < -a/2$$
  

$$II: \quad -a/2 \le z \le a/2$$
  

$$III: \quad z > a/2$$
(8)

- (a) (4 pts) Write down the correct Dirac equation for the motion of this particle.
- (b) (4 pts) Show that the equation found in (a) is invariant under parity transformation, i.e. the transformation  $z \rightarrow -z$
- (c) (6 pts) Solve the equation for the ground state and wave function of the trapped particle. Take the limit  $V_0 \to \infty$ , and sketch the solution for this case. Comment on any interesting feature which this solution might show.